

A THEOREM ON INTEGRAL EQUATION

By V. D. KORANNE

1.—The object of this paper is to establish a theorem on functions of two variables which under suitable conditions satisfy the integral equation

$$f(x, y) =$$

$$x^{\lambda_1} \cdot y^{\lambda_2} \int_0^{\infty} \int_0^{\infty} e^{a_1 x - b_1 u + a_2 y - b_2 v} J_{\lambda_1}^{M_1}(x^{M_1} u) J_{\lambda_2}^{M_2}(y^{M_2} v) f(u, v) du dv$$

where $J_{\lambda}^M(x)$ is Bessel - Maitland Function [1] defined by

$$J_{\lambda}^M(x) = \sum_{\alpha=0}^{\infty} \frac{(-x)^{\alpha}}{\alpha! \Gamma(1 + \lambda + M\alpha)} \quad ; \quad (M > 0).$$

The transform involving above function as Kernel has been studied by R. Kumar [2].

It is interesting to note that this theorem leads to (i) a general expression for functions which are R_M, ν i.e. which satisfy the integral equation

$$f(x, y) =$$

$$\int_0^{\infty} \int_0^{\infty} \sqrt{xytz} J_M(xt) J_{\nu}(yz) f(t, z) dt dz, \quad R(M, \nu) > -1$$

and (ii) a generalised result in two variables analogous to well known Tricomi's Theorem [3].

2.—*Theorem:* The function $f(x, y)$ determined by

$$f(x, y) = \frac{1}{(2\pi i)^2} \int_{\nu_1 - i\infty}^{\nu_1 + i\infty} \int_{\nu_2 - i\infty}^{\nu_2 + i\infty} e^{px + qy} \varnothing(p, q) dpdq \quad (2.1)$$

satisfies the integral equation

$$f(x, y) = x^{\lambda_1} y^{\lambda_2} \int_0^\infty \int_0^\infty e^{a_1 x - b_1 u + a_2 y - b_2 v} J_{\lambda_1}^{M_1}(x^{M_1} u) J_{\lambda_2}^{M_2}(y^{M_2} v) f(u, v) dudv$$

provided the functional relation (2.2)

$$\varnothing(p, q) = \frac{\varnothing\{b_1 + (p - a_1)^{-M_1}, b_2 + (q - a_2)^{-M_2}\}}{(p - a_1)^{1+\lambda_1} (q - a_2)^{1+\lambda_2}} \quad (2.3)$$

holds.

This theorem is valid for

$$(i) \quad f(u, v) = 0 \left(u^{-1+\epsilon_1}, v^{-1+\epsilon_2} \right)$$

when u and v are small and $(\epsilon_1, \epsilon_2) > 0$.

$$(ii) \quad f(u, v) = 0 \left(u^{\beta_1} \cdot e^{S_1 u}, v^{\beta_2} \cdot e^{S_2 v} \right)$$

when u and v are large.

$$(iii) \quad p \equiv \gamma_1 + i\eta_1, q = \gamma_2 + i\eta_2, \\ \gamma_1 > c_1, \gamma_2 > c_2; (M_1, M_2) > 0;$$

$$R(\lambda_1, \lambda_2) > -1; R(p - S_1) > 0, \\ R(p - S_2) > 0, R(b - S_1) > 0$$

$$R(b_2 - S_2) > 0; R(p - a_1) > 0, R(q - a_2) > 0;$$

$$R(c_1 - S_1) > 0, R(c_2 - S_2) > 0.$$

If $R(S_1 - b_1) = R(S_2 - b_2) = 0$, the theorem holds for

$(M_1, M_2) < 1$. For $M_1 = M_2 = 1$, the additional conditions

$$R(\lambda_1 - 2\beta_1) > 0, R(\lambda_2 - 2\beta_2) > 0, \text{ should be satisfied.}$$

Proof: Let $\phi(p, q)$ be Laplace transform of $f(x, y)$ so that

$$\phi(p, q) = \int_0^{\infty} \int_0^{\infty} e^{-px-ky} f(x, y) dx dy;$$

$$R(p - S_1) > 0, R(q - S_2) > 0.$$

Then if $f(x, y)$ satisfied (2.2) we have

$$\phi(p, q) = \int_0^{\infty} \int_0^{\infty} x^{\lambda_1} y^{\lambda_2} e^{-px-ky} dx dy$$

$$\int_0^{\infty} \int_0^{\infty} e^{+a_1x - b_2u + a_2y - b_2v} J_{\lambda_1}^{M_1}(x^{M_1} u) J_{\lambda_2}^{M_2}(y^{M_2} v) f(u, v) dudv$$

$$= \int_0^{\infty} \int_0^{\infty} e^{-b_1u - b_2v} f(u, v) dudv$$

$$\int_0^{\infty} \int_0^{\infty} x^{\lambda_1} y^{\lambda_2} e^{-(p-a_1)x} e^{-(q-a_2)y} J_{\lambda_1}^{M_1}(x^{M_1} u) J_{\lambda_2}^{M_2}(y^{M_2} v) dx dy$$

.. (I)

$$\begin{aligned}
&= \int_0^{\infty} \int_0^{\infty} e^{-b_1 u - b_2 v} f(u, v) \, du \, dv \times \sum_{\alpha_1}^{\infty} \frac{(-u)^{\alpha_1}}{\alpha_1 j \Gamma(1 + \lambda_1 + M_1 \alpha_1)} \times \\
&\sum_{\alpha_2}^{\infty} \frac{(-v)^{\alpha_2}}{\alpha_2 j \Gamma(1 + \lambda_2 + M_2 \alpha_2)} \times \\
&\int_0^{\infty} \int_0^{\infty} x^{\lambda_1 + M_1 \alpha_1} y^{\lambda_2 + M_2 \alpha_2} e^{-(p-a_1)x} e^{-(q-a_2)y} \, dx \, dy \quad \dots (I_2) \\
&= \int_0^{\infty} \int_0^{\infty} e^{-b_1 u - b_2 v} f(u, v) \, du \, dv \times \sum_{\alpha_1}^{\infty} \frac{(-u)^{\alpha_1}}{\alpha_1 j (p-a_1)^{1 + \lambda_1 + M_1 \alpha_1}} \cdot \\
&\sum_{\alpha_2}^{\infty} \frac{(-v)^{\alpha_2}}{\alpha_2 j (q-a_2)^{1 + \lambda_2 + M_2 \alpha_2}} \cdot \\
&x \, du \, dv \\
&= \frac{1}{(p-a_1)^{\lambda_1+1} \cdot (q-a_2)^{\lambda_2+1}} \times \\
&\int_0^{\infty} \int_0^{\infty} e^{-b_1 u - (p-a_1)^{-M_1} \cdot u - b_2 v - (q-a_2)^{-M_2} v} f(u, v) \, du \, dv \\
&= \frac{\emptyset \{b_1 + (p-a_1)^{-M_1} \cdot b_2 + (q-a_2)^{-M_2}\}}{(p-a_1)^{\lambda_1+1} \cdot (q-a_2)^{\lambda_2+2}}
\end{aligned}$$

and then by [4]

$$f(x, y) = \frac{1}{(2\pi i)^2} \int_{V_1-i\infty}^{V_1+i\infty} \int_{V_2-i\infty}^{V_2+i\infty} e^{px+qy} \varnothing(p, q) dp dq$$

provided $p \equiv$

$$\begin{aligned} & \sqrt{1+i\eta_1}, q \equiv \sqrt{2+i\eta_2}, \\ & \sqrt{1}>c_1, \sqrt{2}>c_2, \\ & R(c_1-s_1) > 0, \text{ and } R(c_2-s_2) > 0. \end{aligned}$$

It only remains to justify the inversions in the steps I_1 and I_2 .

We know that

$$\text{as } x \rightarrow \infty, J_\lambda^M(x) = 0 \quad \left[x^{-k(a+1/2)} \exp. \{ (Mx)^k \frac{\text{Cos}(\eta k)}{Mk} \} \right]$$

$$k = \frac{1}{M+1} \text{ and as } x \rightarrow 0 J_\lambda^M(x) = 0 \quad (1)$$

Therefore both integrals converge absolutely by virtue of De La Valle Poussin's Theorem [5], the inversion in the step I_1 is justified.

La inversion in the step I_2 is valid since

(i) each term

$$T_\alpha(x) \equiv \frac{(-x^M u)^\alpha}{\alpha j \Gamma(1+\lambda+M\alpha)}$$

of the series is continuous.

(ii) $\Sigma |T_\alpha(x)|$ is uniformly convergent in arbitrary large interval $(0, \alpha)$

and (iii)

$$\sum_{\alpha=0}^{\infty} \int_0^{\infty} \frac{\lambda}{x} e^{-(p-a)x} |T_\alpha(x)| dx$$

$\alpha=0 \quad 0$

converges to

$$(p-a)^{-\lambda-1} \exp. \{ (p-a)^{-M} u \} .$$

Example:

If we take

$$\phi(p, q) = \frac{\Gamma\nu_1}{(p-a_1)^{\nu_1}} \times \frac{\Gamma\nu_2}{(p-a_2)^{\nu_2}} ; R(\nu_1, \nu_2) > 0$$

and

$$b_1 = a_1, b_2 = a_2$$

then the functional relation (2.3) is satisfied provided.

$$M_1 = \frac{\lambda_1 + 1 - \nu_1}{\nu_1}, \quad M_2 = \frac{\lambda_2 + 1 - \nu_2}{\nu_2}$$

Therefore

$$f(x, y) = \frac{1}{(2\pi i)^2} \int_{r_1 - i\infty}^{r_2 + i\infty} \int_{r_2 - i\infty}^{r_2 + i\infty} e^{px+qy} \frac{\Gamma\nu_1}{(p-a_1)^{\nu_1}} \cdot \frac{\Gamma\nu_2}{(q-a_2)^{\nu_2}} dpdq$$

$$= x^{\nu_1-1} e^{a_1 x} \cdot y^{\nu_2-1} e^{a_2 y},$$

satisfies the integral equation (2.2) with $b_1 = a_1, b_2 = a_2$

and

$$M_1 = \frac{\lambda_1 + 1 - \nu_1}{\nu_1}, \quad M_2 = \frac{\lambda_2 + 1 - \nu_2}{\nu_2}$$

provided

$$p \equiv r_1 + i\eta_1, q \equiv r_2 + i\eta_2; r_1 > c_1, r_2 > c_2;$$

$$M_1 > 0, M_2 > 0,$$

$R(p-a_1, q-a_2) > 0$; $R(c_1-a_1, c_2-a_2) > 0$;
 $R(2\nu_1, 2\nu_2) > 0$, $R(\lambda_1+1, \lambda_2+1) > 0$.

3.—Corollary. If $f(x, y)$ is the function determined by

$$f(x, y) = \frac{1}{(2\pi i)^2} \int_{r_1-i\infty}^{r_1+i\infty} \int_{r_2-i\infty}^{r_2+i\infty} a^{+px+ay} \varnothing(p, q) dpdq \quad (3.1)$$

and $\varnothing(p, q)$ satisfies the functional relation

$$\varnothing(p, q) = \frac{\varnothing\{a_1+(p_1-a_1)^{-1} \quad a_2+(q-a_2)^{-1}\}}{(p-a_1)^{\nu_1+1} (q-a_2)^{\nu_2+2}}$$

then

$$e^{-a_1 \frac{x^2}{2} - a_2 \frac{y^2}{2}} \cdot x^{-\nu_1+1/2} y^{-\nu_2+1/2} f\left(\frac{x^2}{2}, \frac{y^2}{2}\right)$$

is $R\nu_1, \nu_2$

provided

(i) $f(x, y) = 0 \left(x^{\frac{-1+\epsilon_1}{2}}, y^{\frac{-1+\epsilon_2}{2}}\right)$
when x and y are small ($\epsilon_1, \epsilon_2 > 0$).

(ii) $f(x, y) = 0 \left(x^{\frac{\beta_1 s_1 x}{e}}, y^{\frac{\beta_2 s_2 y}{e}}\right)$
when x and y are large.

(iii) $p \equiv r_1 + i\eta_1, q \equiv r_2 + i\eta_2; r_1 > c_1, r_2 > c_2; R(p, q) > 0$;
 $R(\nu_1 - \nu_2) > -1; R(p_1 - s_1, q - s_2) > 0; R(p - a_1, q - a_2) > 0$;
 $R(s_1 - a_1, s_2 - a_2) < 0; R(c_1 - s_1, c_2 - s_2) > 0$.

If $R(s_1 - a_1, s_2 - a_2) = 0$, the additional condition

$R(\nu_1 - 2\beta_1, \nu_2 - 2\beta_2) > -1/2$ must be satisfied. For on putting

$M_1=1, M_2=1, a_1=b_2$ and $\lambda_1=\nu_1, \lambda_2=\nu_2$ (2.3) reducesto (3.2) and gives

$M_1=1, M_2=1, a_1=b_2, a_2=b_2$ and $\lambda_1=\nu_1, \lambda_2=\nu_2$ (2.3)

reducesto (3.2) and gives

$$f(x, y) = x^{\nu_1} y^{\nu_2} \int_0^{\infty} \int_0^{\infty} e^{a_1x - a_1u + a_2y - a_2v} J_{\nu_1}^1(xu) J_{\nu_2}^1(yv) f(u, v) dudv$$

or

$$x^{\frac{-\nu_1}{2}} y^{\frac{-\nu_2}{2}} e^{a_1x - a_2y} f(x, y)$$

$$\int_0^{\infty} \int_0^{\infty} J_{\nu_1}(2\sqrt{xu}) J_{\nu_2}(2\sqrt{yv}) u^{-\nu_1/2} v^{-\nu_2/2} e^{-a_1u - a_2v} f(u, v) dudv$$

Replacing x by $x^2/2$, y by $y^2/2$, u by $u^2/2$ and v by $v^2/2$ we get

$$x^{-\nu_1 + 1/2} y^{-\nu_2 + 1/2} e^{-a_1x^2/2 - a_2y^2/2} f\left(\frac{x^2}{2}, \frac{y^2}{2}\right)$$

$$= \int_0^{\infty} \int_0^{\infty} \sqrt{xyuv} J_{\nu_1}(xu) J_{\nu_2}(yv)$$

$$u^{-\nu_1 + 1/2} v^{-\nu_2 + 1/2} e^{+a_1\frac{u^2}{2} + a_2\frac{v^2}{2}} x f\left(\frac{u^2}{2}, \frac{v^2}{2}\right) dudv$$

i.e.

$$x^{-\nu_1 + 1/2} y^{-\nu_2 + 1/2} e^{-a_2\frac{x^2}{2} - a_1\frac{y^2}{2}} f\left(\frac{x^2}{2}, \frac{y^2}{2}\right)$$

is R_{ν_1, ν_2}

On putting $a_1 = 0$, $a_2 = 0$, we get.

if $f(x, y)$ is the function determined by

$$f(x, y) = \frac{1}{(2\pi i)^2} \int_{r_1 - i\infty}^{r_1 + i\infty} \int_{r_2 - i\infty}^{r_2 + i\infty} e^{px+qy} \varnothing(p, q) dpdq$$

and $\varnothing(p, q)$ satisfies the functional relation

$$\varnothing(p, q) = p^{-\nu_1-1} q^{-\nu_2-1} \varnothing\left(\frac{1}{p}, \frac{1}{q}\right)$$

then under suitable conditions

$$x^{-\nu_1+1/2} y^{-\nu_2+1/2} f\left(\frac{x^2}{2}, \frac{y^2}{2}\right)$$

is R_{ν_1, ν_2} .

4.—From § 2, the following theorem is evident

Si $\varnothing(p, q) \doteq f(u, v)$

Then

$$\frac{1}{(p-a_1)^{\lambda_1+1} (p-a_2)^{\lambda_2+1}} \varnothing\{b_1 + (p-a_1)^{-M_1}, b_2 + (q-a_2)^{-M_2}\}$$

$$\doteq \doteq x^{\lambda_1} y^{\lambda_2} \int_0^\infty \int_0^\infty e^{a_1x - b_1u + a_2y - b_2v} J_{\lambda_1}^{M_1}(x^{M_1}u) J_{\lambda_2}^{M_2}(x^{M_2}u) f(u, v) dudv$$

provided

(i) $(u, v) = O\left(x^{-1+\epsilon_1}, y^{-1+\epsilon_2}\right)$
when u and v are small

- (ii) $f(u, v) = 0$ ($u \sim e^{-\beta_1 s_1 u}$, $v \sim e^{-\beta_2 s_2 v}$)
when u and v are large
- (iii) $R(\lambda_1, \lambda_2) > -1$; $R(p-a_1, q-a_2) > 0$; $R(s_1-b_1, s_2-b_2) < 0$,
 $(M_1, M_2) > 0$; $R(p-s_1, q-s_2) > 0$.

If $f(u, v) = 0$ ($u \sim e^{-\beta_1 s_1 u}$, $v \sim e^{-\beta_2 s_2 v}$) and $R(s_1-b_1, s_2-b_2) = 0$ then theorem holds for $(M_1, M_2) < 1$.

It also holds for $M_1 = 1, M_2 = 1$ provided additional condition $R(\lambda_1-2\beta_1, \lambda_2-2\beta_2) > -1/2$ is satisfied. In particular if $a_1=a_2=b_1=b_2=0$ and $M_1=M_2=1$, we obtain Tricomi's theorem in two variables.

If $\mathcal{O}(p, q) \doteq \doteq f(u, v)$, then under suitable conditions

$$\int_0^\infty \int_0^\infty \left(\frac{x}{u}\right)^{\lambda_1/2} \left(\frac{y}{v}\right)^{\lambda_2/2} J_{\lambda_1}(2\sqrt{xu}) J_{\lambda_2}(2\sqrt{yv}) f(u, v) dudv$$

5.—A number of integrals can be evaluated with the help of this theorem.

As an illustration let us take

$$\mathcal{O}(p, q) = \frac{1}{p^{\nu_1+1} q^{\nu_2+1}}$$

So that

$$f(x, y) = \frac{x^{\nu_1} \cdot y^{\nu_2}}{\Gamma(\nu_1+1) \Gamma(\nu_2+1)}; R(\nu_1, \nu_2) > -1$$

and

$$\frac{1}{p^{\lambda_1+1} q^{\lambda_2+1}} \mathcal{O}\{1+p^{-M_1}, 1+q^{-M_2}\}$$

$$= \frac{M_1 (\nu_1 + 1) - \lambda_1 - 1}{p} \frac{M_2 (\nu_2 + 1) - \lambda_2 - 1}{q}$$

$$= \frac{M_1 \nu_1 + 1}{(1+p)} \frac{M_2 \nu_2 + 1}{(1+q)}$$

Case I. Cuando $M_1=M_2=3$; $\nu_1=\nu_2=-\frac{1}{2}$ and $\lambda_1=\lambda_2=\frac{1}{2}$, we have

$$\frac{1}{p^{3/2} q^{3/2}} \mathcal{O} \{1+p^{-3}, 1+q^{-3}\}$$

$$= \frac{1}{(1+p^3)^{1/2} (1+q^3)^{1/2}}$$

$$\doteq \doteq \frac{4}{9} \pi \sqrt{xy} J_{\frac{1}{6}, -\frac{1}{6}}(x) \cdot$$

$$J_{\frac{1}{6}, -\frac{1}{6}}(y)$$

where

$$J_M y \nu (x) = \frac{x^{M+\nu}}{\Gamma(M+1) \Gamma(\nu+1)} x \mathcal{O} F_2 (M+1, \nu+1; -\frac{x^3}{27}).$$

Therefore

$$\int_0^\infty \int_0^\infty u^{-1/2} v^{-1/2} J_{\frac{1}{6}}^3 (x^3 u) J_{\frac{1}{6}}^3 (y^3 v) e^{-u-v} dudv$$

$$= \frac{4\pi^2}{9} J_{\frac{1}{6}, -\frac{1}{6}}(x) \cdot J_{\frac{1}{6}, -\frac{1}{6}}(y)$$

Case II. When $M_1=M_2=3$; $\nu_1=\nu_2=-\frac{1}{2}$ and $\lambda_1=\lambda_2=-\frac{1}{2}$

we have

$$p^{-1/2} q^{-1/2} \oslash (1+p^{-3}, 1+q^{-3}) = \frac{pq}{(p^3+1)^{1/2} (q^3+1)^{1/2}}$$

$$\doteq \doteq \frac{4}{9} \pi \sqrt{xy} J_{-1/6, -5/6}(x) \cdot J_{-1/6, -5/6}(y)$$

and therefore

$$\int_0^{\infty} \int_0^{\infty} u^{-1/2} v^{1/2} e^{-u-v} J_{-1/2}^3(x^3 u) J_{-1/2}^3(y^3 v) dudv$$

$$= \frac{4\pi^2}{9} xy J_{1/6, -5/6}(x) \cdot J_{-1/6, -5/6}(y)$$

Case III. When $M_1=M_2=2$; $\nu_1=\nu_2=0$; $\lambda_1=\lambda_2=1$

$$\frac{1}{p^2 q^2} \oslash \left(1 + \frac{1}{p^2}, 1 + \frac{1}{q^2}\right) = \frac{1}{(1+p^2)(1+q^2)} \doteq \doteq \sin x \sin y$$

and therefore

$$\int_0^{\infty} \int_0^{\infty} e^{-u-v} J_0^2(x^2 u) J_0^2(y^2 v) du dv = \frac{\sin x \sin y}{xy}$$

Case IV. When $M_1=M_2=2$; $\nu_1=\nu_2=0$; $\lambda_1=\lambda_2=0$

$$\frac{1}{pq} \oslash \left(1 + \frac{1}{p^2}, 1 + \frac{1}{q^2}\right) = \frac{1}{(1+p^2)(1+q^2)} \doteq \doteq \cos x \cdot \cos y$$

and therefore

$$\int_0^{\infty} \int_0^{\infty} e^{-u-v} J_0^2(x^2 u) J_0^2(y^2 v) dudv = \cos x \cdot \cos y$$

Case V. When $M_1=M_2=2$; $\nu_1=\nu_2=-\frac{1}{2}$; $\lambda_1=\lambda_2=0$

$$\frac{1}{pq} \mathcal{O} \left(1 + \frac{1}{p^2}, 1 + \frac{1}{q^2} \right) = \frac{1}{(1+p^2)^{1/2} (1+q^2)^{1/2}} \doteq \doteq J_0(x) \cdot J_0(y)$$

and therefore

$$\int_0^\infty \int_0^\infty e^{-u-v} J_0^2(x^2 u) J_0^2(y^2 v) du dv$$

$$= \pi J_0(x) \cdot J_0(y).$$

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R E F E R E N C E S

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